

# Improvement and protection of niobium surface superconductivity by atomic layer deposition and heat treatment

FERMILAB-PUB-08-543-TD

T. Proslir,<sup>1,2,a)</sup> J. Zasadzinski,<sup>1</sup> J. Moore,<sup>3</sup> M. Pellin,<sup>2</sup> J. Elam,<sup>4</sup> L. Cooley,<sup>5</sup> C. Antoine,<sup>6</sup> J. Norem,<sup>3</sup> and K. E. Gray<sup>2</sup>

<sup>1</sup>Physics Department, Illinois Institute of Technology, Chicago, Illinois 60616, USA

<sup>2</sup>Materials Science Division, Argonne National Laboratory, Argonne, Illinois 60493, USA

<sup>3</sup>High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois 60493, USA

<sup>4</sup>Energy Systems Division, Argonne National Laboratory, Argonne, Illinois 60493, USA

<sup>5</sup>Technical Division, Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA

<sup>6</sup>Commissariat à l'énergie Atomique, Centre de Saclay, F-91191, Gif-Sur-Yvette, France

(Received 13 August 2008; accepted 16 September 2008; published online 11 November 2008)

A method to treat the surface of Nb is described, which potentially can improve the performance of superconducting rf cavities. We present tunneling and x-ray photoemission spectroscopy measurements at the surface of cavity-grade niobium samples coated with a 3 nm alumina overlayer deposited by atomic layer deposition. The coated samples baked in ultrahigh vacuum at low temperature degraded superconducting surface. However, at temperatures above 450 °C, the tunneling conductance curves show significant improvements in the superconducting density of states compared with untreated surfaces.

□

For several decades, the performances of accelerating niobium superconducting radio frequency (SRF) cavities have been continually improved to reach now a reproducible maximum accelerating field of  $\sim 30$  MV/m and quality factor  $Q$  above  $10^{10}$ . However, unresolved mysteries remain<sup>1</sup> and prevent cavities from reaching the intrinsic Nb limits believed to be 55 MV/m. An SRF cavity is subject to an accelerating rf field  $\mathbf{E}(\omega)$  that induces an oscillating magnetic field  $\mathbf{B}(\omega)$  on the inner walls of the cavity. Penetration of the latter implies superconducting screening currents on a length scale  $\lambda_{\text{Nb}}$ , the niobium penetration depth ( $\sim 45$  nm at  $T=2$  K). The dissipation mechanisms are confined to a region of depth  $\lambda_{\text{Nb}}$  over which supercurrents interact strongly with the surface oxides. Up to recently,<sup>2</sup> a detailed picture of these strong interactions and, in particular, the mechanisms by which oxygen influences the performance of niobium cavities was still missing.

The uncontrolled complex sets of oxides growing on the surface of air-exposed niobium lead to well-known problematic effects such as weakened superconductivity, proximity effects, or pair breaking phenomena. For niobium-based devices,<sup>3</sup> submicron Josephson tunnel junctions or single electron transistors of these problems have been cured by capping a cleaned niobium surface with alumina. However, at present there are no realistic ways to remove the oxides from the inner wall of niobium cavities and to deposit a uniformly thin protective layer on a complex-shaped cavity. We propose an original approach to overcome these problems: the combination of the atomic layer deposition (ALD) technique with a subsequent annealing in ultrahigh vacuum (UHV). We probe the occurrence of surface oxides with x-ray photoemission spectroscopy (XPS) and the near surface superconductivity by tunneling spectroscopy.

Tunneling spectroscopy and XPS measurements were done on air-exposed cavity-grade niobium coupons capped

with a protective alumina layer deposited by ALD.<sup>4</sup> ALD is a self-limiting, sequential surface chemistry that has the unique ability to achieve uniform atomic scale deposition control on arbitrary complex-shaped substrates.<sup>5</sup> The pinhole-free and dense structure of ALD made films make  $\text{Al}_2\text{O}_3$  an ideal protective and diffusion barrier. Pieces of monocrystalline (110) and polycrystalline Nb were cut from larger sheets used to construct SRF cavities. These pieces were electropolished, cleaned with de-ionized purified water, and dried in air in a manner similar to that done on cavities. They were later introduced into the ALD reactor and coated with 3 nm of  $\text{Al}_2\text{O}_3$  at a temperature of 120 °C under a flow of ultrapure nitrogen gas. A detailed description of the deposition process can be found in Ref. 6. The coated samples were baked in UHV at temperature ranging from 250 to 500 °C and the XPS spectrum of the Nb 3d core level was measured simultaneously. After each thermal treatment, the samples were transferred in air to the point contact apparatus and the surface superconducting density of states (DOS) was probed by tunneling spectroscopy at a temperature of 1.7 K as described elsewhere.<sup>7</sup>

The evolution of the oxide's composition and the superconducting DOS at the surface of a polycrystalline sample as a function of the baking temperature is shown in Fig. 1. For annealing temperature up to 380 °C, the tunneling spectroscopy [Fig. 1(a)] reveals degraded superconducting features as compared to unbaked samples (in blue) with a zero bias conductance (ZBC) as high as 65% of the normal conductance (measured at  $V=12$  meV). Importantly, for an anneal at 450 °C for 24 h, the tunneling conductance becomes much sharper and the ZBC improves by a factor of 5 down to  $\sim 5\%$ . In parallel, the XPS spectrum [Fig. 1(b)] that probes an  $\sim 3.5$  nm depth on the surface shows a reduction in  $\text{Nb}_2\text{O}_5$  to suboxides such as NbO and  $\text{NbO}_2$  below 450 °C (in green and orange), whereas at 450 °C 24 h the XPS unveils the presence of NbO only with sharper metallic Nb<sup>0</sup> peaks. The same reduction in the ZBC down to  $\sim 5\%$  is found reproducibly on Nb single crystals [Fig. 1(c)] baked at

<sup>a)</sup>Electronic mail: proslir@anl.gov.

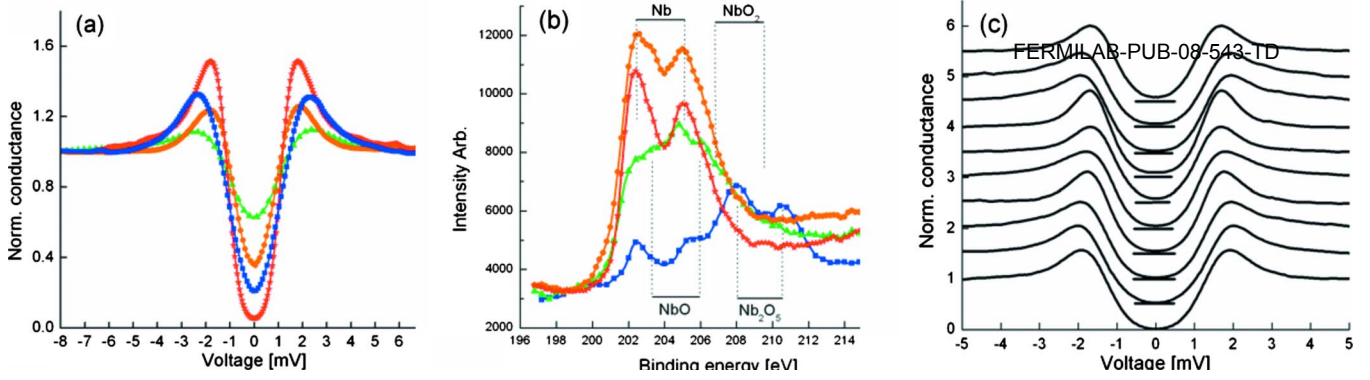


FIG. 1. (Color online) (a) Conductance curves measured at 1.7 K on a polycrystalline Nb sample coated with 3 nm of  $\text{Al}_2\text{O}_3$ . (b) Corresponding XPS spectrum of the Nb 3d core level. Each color and symbol correspond to one temperature: 30 °C in blue squares, 220 °C in green triangles, 380 °C in orange circles, and 450 °C in red stars. (c) Ten conductance curves measured on a coated Nb single crystal (110) baked at 500 °C for 24 h. The curves have been shifted by 0.5 for more clarity and the corresponding zero conductance references for each curve is represented by short lines around the Fermi level ( $V=0$  mV).

500 °C for 24 h in UHV. We interpret the abrupt decrease in the ZBC above 450 °C as a consequence of the oxygen diffusion from Nb oxides into the bulk, leading to a cleaner interface and to a reduced inelastic scattering at the surface.

In the present study, we observe a shift in the quasiparticle peaks to voltages above the bulk Nb energy gap of 1.55 meV, suggesting an inelastic scattering mechanism. However, while baking the situation is more complex and well-known deleterious effects (massive injection of oxygen into Nb) do occur as well as other elastic processes that could lead to filling of the gap. To extract the proportion of inelastic and elastic scattering mechanisms is beyond the scope of this paper and perhaps impossible. As a consequence, we think of  $\Gamma$  as a broadening parameter that gives a quantitative estimate of the deviation of the measured superconducting DOS from an ideal one. The evolution of  $\Delta$  and  $\Gamma$  as a function of the baking temperature is summarized in Fig. 2(a).

It is worth noting that for an unbaked coated sample, the fitting parameters are the same as for uncoated Nb samples, i.e., the deposition of an alumina overlayer by ALD does not modify the superconducting parameters at the surface of Nb. However, for baking temperatures of 250 °C 2 h and 380 °C 24 h, the superconducting DOS is strongly impaired, exhibiting reduced gap values of 1.35 and 1.3 meV, respectively. The evolution of the gap coincides with the increasing peak

intensity of the total niobium oxides extracted from the XPS fits<sup>8</sup> [Fig. 2(b)], mostly composed of metallic NbO as the  $\text{Nb}_2\text{O}_5$  is reduced. This result is in agreement with previous XPS studies<sup>9–12</sup> that indeed revealed a thickening of metallic NbO at the interface with the niobium. The smaller measured energy gap may thus be a result of a proximity effect between the Nb and a thicker NbO. The corresponding variations in the inelastic scattering value  $\Gamma$ , jumping from 0.6 meV at 30 °C to 1 meV at 250 °C 2 h and decreasing back to 0.48 meV at 380 °C 24 h, do not seem to correlate with the oxide composition. This can be explained, however, by the competition between the strong oxide dissociation above 200 °C (Ref. 13) and the diffusion of oxygen into Nb. At 250 °C for 2 h, the diffusion length  $L_{\text{diff}}$  calculated from the diffusion equation of oxygen into the bulk niobium is 120 nm  $\sim 2\lambda_{\text{Nb}}$ , whereas above 380 °C for longer time, 24 h, the diffusion is faster ( $L_{\text{diff}} \geq 1$  mm) and oxygen is driven far into the bulk. Consequently the  $\Gamma$  value decreases. Interestingly, the evolution of the superconducting parameters correlates well with cavity performance after similar baking treatment.<sup>14</sup>

Finally, for the highest baking temperature, >450 °C for 24 h, the total intensity of niobium oxide peaks composed of only NbO decreases drastically in agreement with Ref. 10 and with the idea that oxygen diffuses far into the bulk. This

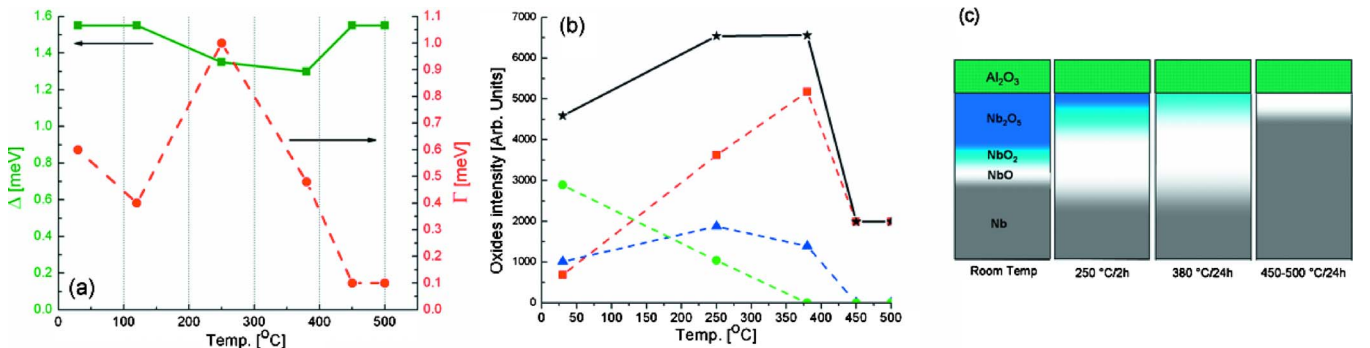


FIG. 2. (Color online) (a) Evolution of the superconducting gap  $\Delta$  (green squares) and inelastic scattering parameter  $\Gamma$  (red circles) as a function of the baking temperature. (b) Plot of the intensity of Nb oxide peaks extracted from XPS fits, the solid black line represents the total intensity of oxide peaks, in red squares are the NbO peak intensity, blue triangles the  $\text{NbO}_2$  and in green circles the  $\text{Nb}_2\text{O}_5$  ones. (c) Schematic evolution of the oxide composition at the surface of Nb. The various shades account for the gradual transition between two different oxides composition, as the corresponding interfaces are rather blurred and continuous than abrupt and well ordered.<sup>1</sup>

result in a striking improvement in the superconducting properties:  $\Gamma=0.1$  meV and  $\Delta$  recovers the bulk Nb gap value of 1.55 meV. The saturation of the NbO peak intensity above 450 °C may be an indication of a more stable and ordered NbO phase, as proposed by Hellwig *et al.*,<sup>15</sup> with a sharper NbO–Nb interface that would also contribute to improve the proximity effect of the superconducting DOS. Higher temperature anneals will be carried out on coated niobium coupons to search for further improvements in the superconducting DOS, aiming at surface properties of an ideal Nb superconductor with no inelastic scattering and optimum Nb cavity performance.

In conclusion, the overall evolution of the superconducting parameters  $\Gamma$  and  $\Delta$ , together with the oxide composition as a function of the baking temperature, can be explained by a competition between Nb-oxide dissociation and oxygen diffusion into the bulk. Such a model has been suggested for low temperature (<150 °C) anneals by Ciovati,<sup>13</sup> but it has not been proven experimentally for high baking temperature. The close correlation of the results in Fig. 2(a) with cavity performance for the same temperature and baking time points to tunneling spectroscopy as a relevant tool to explain and predict cavity performance and inexpensively test various surface treatment protocols. A single cell 1.3 GHz cavity has been coated with 10 nm Al<sub>2</sub>O<sub>3</sub> and 3 nm Nb<sub>2</sub>O<sub>5</sub> (to avoid multipacking effect) and rf performance test will be conducted soon.

This work was supported by U.S. DOE under Contract No. DE-AC02-06CH11357.

- <sup>1</sup>J. Halbritter and P. Kneisel, *IEEE Trans. Appl. Supercond.* **11**, 1864 (2001); P. Kneisel and C. Z. Antoine, Proceedings of the Ninth Workshop on RF Superconductivity, edited by F. Krawczyk (Los Alamos National Laboratory, Sante Fe, NM, 2000), pp. 295–303.
- <sup>2</sup>Th. Proslier, J. F. Zasadzinski, J. Moore, L. Cooley, C. Antoine, M. Pellin, J. Norem, and K. Gray, *Appl. Phys. Lett.* **92**, 212505 (2008); A. Gurevich and C. Ciovati, *Phys. Rev. B* **77**, 104501 (2008).
- <sup>3</sup>L. A. Abelson and G. L. Kerber, *Proc. IEEE* **92**, 1515 (2004).
- <sup>4</sup>M. Ritala, M. Leskela, and H. S. Nalwa (Ed.), in *Handbook of Thin Film Materials* (Academic, San Diego, 2001), Vol. 1, p. 103.
- <sup>5</sup>J. W. Elam, D. Routkevitch, P. P. Mardilovich, and S. M. George, *Chem. Mater.* **15**, 3507 (2003).
- <sup>6</sup>J. W. Elam, M. D. Groner, and S. M. George, *Rev. Sci. Instrum.* **73**, 2981 (2002); A. W. Ott, J. W. Klaus, and J. M. George, *Thin Solid Films* **292**, 135 (1997).
- <sup>7</sup>L. Ozyuzer, J. F. Zasadzinski, and K. E. Gray, *Cryogenics* **38**, 911 (1998).
- <sup>8</sup>For spectral analysis we used a curve simulation routine that fits the spectral line shape using Gaussian functions. The ratio of the  $3d_{3/2}$ – $3d_{5/2}$  intensities was set to 1.4 and the difference in energy was kept to 2.74 eV for all oxides and Nb peaks. The atomic concentration of oxygen shows the exact same evolution in temperature as the total niobium oxide peak intensity in Fig. 2(b). The atomic concentration of Nb increases as the temperature is increased with a plateau around 350 °C.
- <sup>9</sup>Q. Ma, P. Ryan, J. W. Freeland, and R. A. Rosenberg, *J. Appl. Phys.* **96**, 7675 (2004).
- <sup>10</sup>M. Delheusy, A. Steirle, N. Kasper, R. Kurta, A. Vlad, H. Dosh, C. Antoine, A. Resta, E. Lundgren, and J. Andersen, *Appl. Phys. Lett.* **92**, 101911 (2008).
- <sup>11</sup>Q. Ma and R. A. Rosenberg, *Appl. Surf. Sci.* **206**, 209 (2003).
- <sup>12</sup>M. Grundner and J. Halbritter, *J. Appl. Phys.* **51**, 397 (1980).
- <sup>13</sup>G. Ciovati, *Appl. Phys. Lett.* **89**, 022507 (2006).
- <sup>14</sup>B. Visentin, J. P. Charrier, A. Aspart, Y. Gasser, J. P. Poupeau, and G. Congretel, Proceedings of the Eighth European Particle Accelerator Conference, Paris, France, June 2002 (unpublished), p. 2292.
- <sup>15</sup>O. Hellwig, H. W. Beker, and H. Zabel, *Phys. Rev. B* **64**, 233404 (2001).